

PHYS 320 ANALYTICAL MECHANICS

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Fourier series:

- Full differential equation of motion:

$$m\ddot{x}(t) + c\dot{x}(t) + kx(t) = F_{ext}(t)$$

- forcing function written as Fourier series:

$$F_{ext}(t) = \sum_{n=1}^{\infty} F_n = \sum_{n=1}^{\infty} b_n \sin(n\omega t)$$

- solution to full diff. eqn. is linear superposition of solutions to the n^{th} differential equation:

$$m\ddot{x}_n(t) + c\dot{x}_n(t) + kx_n(t) = F_n(t)$$

$$x(t) = \sum_{n=1}^{\infty} x_n(t) = \sum_{n=1}^{\infty} A_n \cos(n\omega t - \delta_n)$$

$$\text{where } A_n = \frac{b_n / m}{\left[(\omega_o^2 - n^2 \omega^2)^2 + 4\gamma^2 n^2 \omega^2 \right]^{1/2}} \quad \text{and} \quad \delta_n = \tan^{-1} \left(\frac{2\gamma n \omega}{\omega_o^2 - n^2 \omega^2} \right)$$

Fourier Series

For a function defined on the interval $[-L, L]$, where

$$f(x') = \frac{1}{2} a_0 + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi x'}{L}\right) + \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x'}{L}\right).$$

where

$$\begin{aligned} a_0 &= \frac{1}{L} \int_{-L}^L f(x') dx' \\ a_n &= \frac{1}{L} \int_{-L}^L f(x') \cos\left(\frac{n\pi x'}{L}\right) dx' \\ b_n &= \frac{1}{L} \int_{-L}^L f(x') \sin\left(\frac{n\pi x'}{L}\right) dx'. \end{aligned}$$

Similarly, the function is instead defined on the interval $[0, 2L]$, the above equations simply become

$$\begin{aligned} a_0 &= \frac{1}{L} \int_0^{2L} f(x') dx' \\ a_n &= \frac{1}{L} \int_0^{2L} f(x') \cos\left(\frac{n\pi x'}{L}\right) dx' \\ b_n &= \frac{1}{L} \int_0^{2L} f(x') \sin\left(\frac{n\pi x'}{L}\right) dx'. \end{aligned}$$

Fourier Series

For our context, rewrite in terms of period, T:

$$f(t) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} [a_n \cos(n\omega t) + b_n \sin(n\omega t)]$$

where

$$a_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \cos(n\omega t) dt, \quad n = 0, 1, 2, \dots$$

$$b_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \sin(n\omega t) dt, \quad n = 1, 2, \dots$$

$$T = \frac{2\pi}{\omega}$$

Fourier series: sawtooth driving force

- Solutions are $x(t) = \sum_{n=1}^{\infty} x_n(t) = \sum_{n=1}^{\infty} A_n \cos(n\omega t - \delta_n)$

where $A_n = \frac{b_n / m}{\left[(\omega_o^2 - n^2 \omega^2)^2 + 4\gamma^2 n^2 \omega^2 \right]^{1/2}}, \quad \delta_n = \tan^{-1} \left(\frac{2\gamma n \omega}{\omega_o^2 - n^2 \omega^2} \right)$

$$b_n = \frac{A}{n\pi} (-1)^{n+1}$$

$$b_1 = \frac{A}{\pi}, \quad b_2 = \frac{-A}{2\pi}, \quad b_3 = \frac{A}{3\pi}, \quad \text{etc.}$$

So, we have

$$A_1 = \frac{A / \pi m}{\left[(\omega_o^2 - \omega^2)^2 + 4\gamma^2 \omega^2 \right]^{1/2}}, \quad \delta_1 = \tan^{-1} \left(\frac{2\gamma \omega}{\omega_o^2 - \omega^2} \right)$$

$$A_2 = \frac{-A / 2\pi m}{\left[(\omega_o^2 - 4\omega^2)^2 + 4\gamma^2 (4\omega^2) \right]^{1/2}}, \quad \delta_2 = \tan^{-1} \left(\frac{2\gamma(2\omega)}{\omega_o^2 - 4\omega^2} \right)$$

$$A_3 = \frac{-A / 2\pi m}{\left[(\omega_o^2 - 9\omega^2)^2 + 4\gamma^2 (9\omega^2) \right]^{1/2}}, \quad \delta_3 = \tan^{-1} \left(\frac{2\gamma(3\omega)}{\omega_o^2 - 9\omega^2} \right)$$

et cetera!

if any are near ω_r then get strong response!

